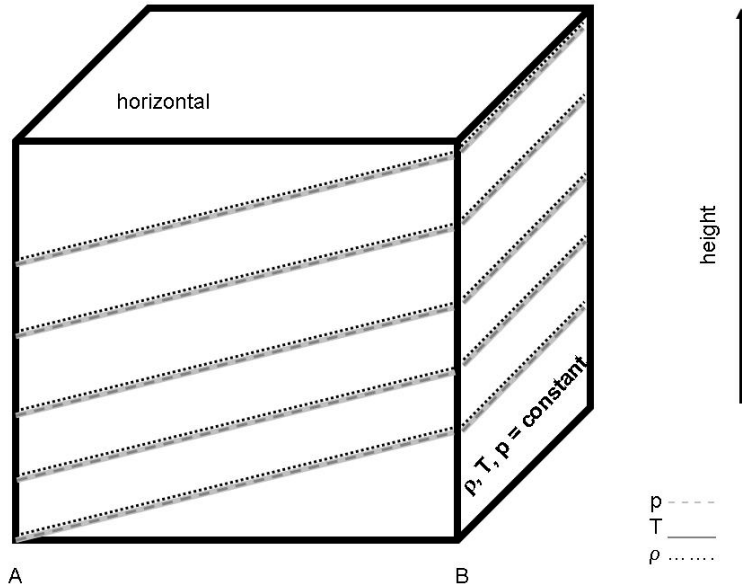


Unit 20

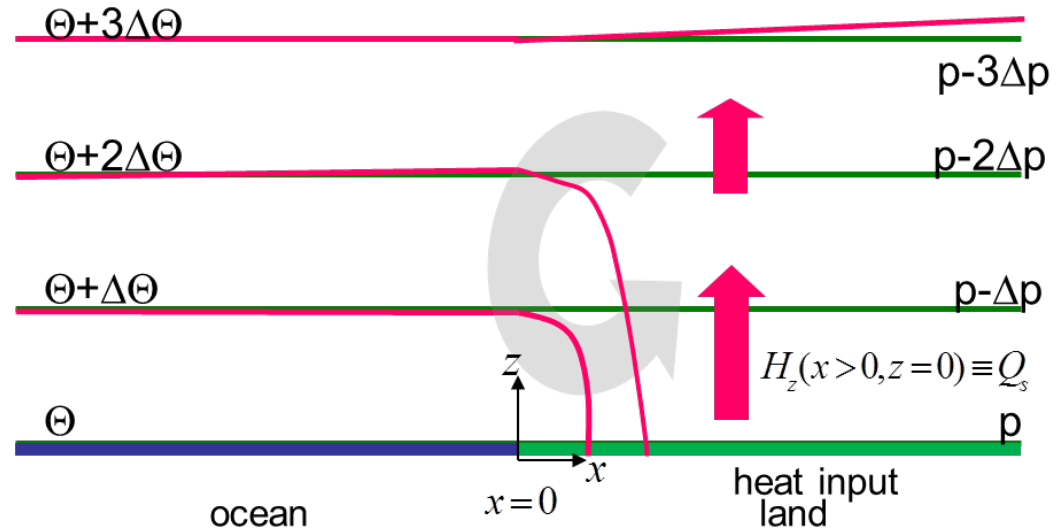
Thermal wind, advection, and primitive equations

Nicole Mölders

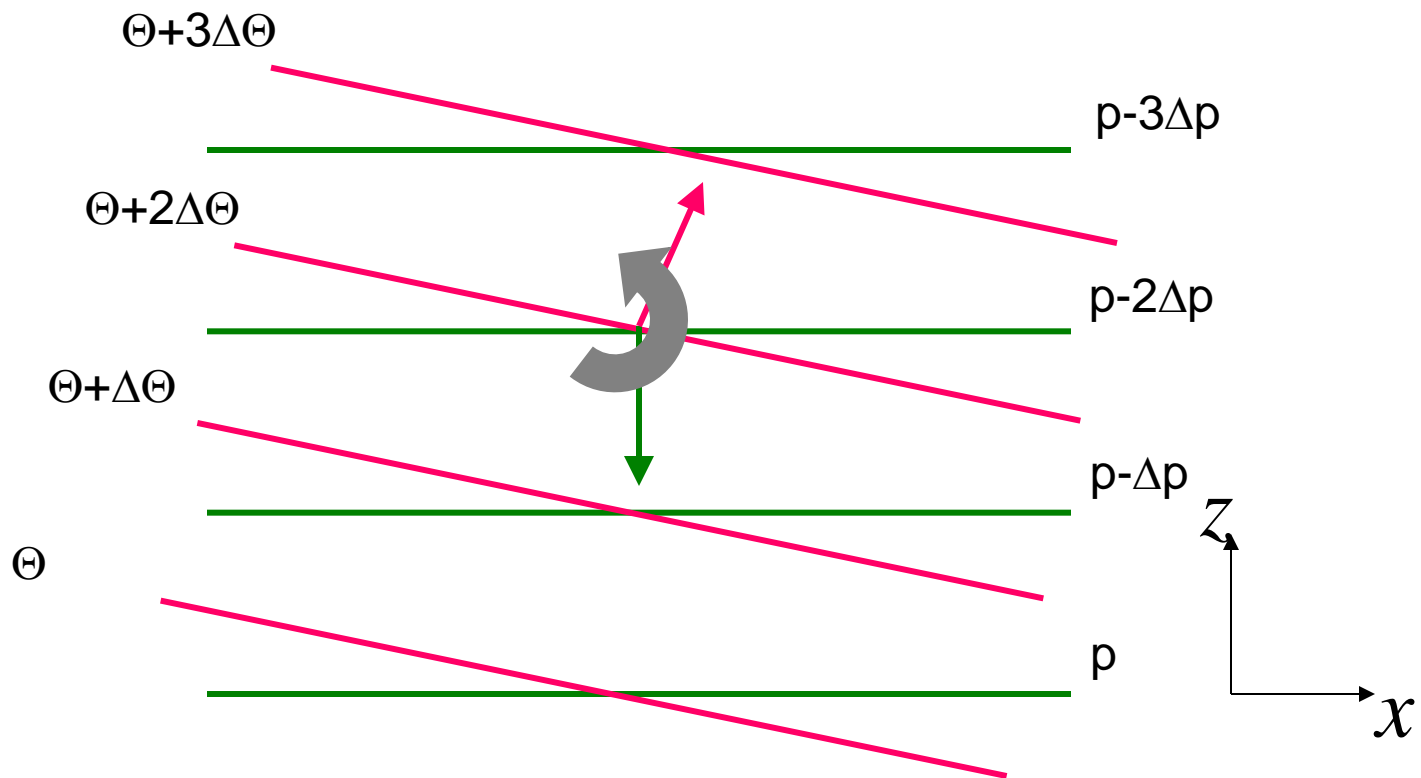
Barotrop vs. barocline



p ---
 T ———
 ρ



Baroclinicity creates vorticity



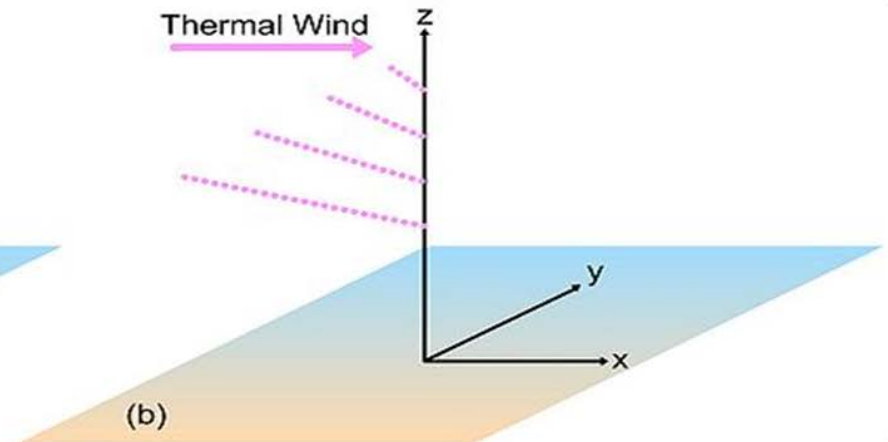
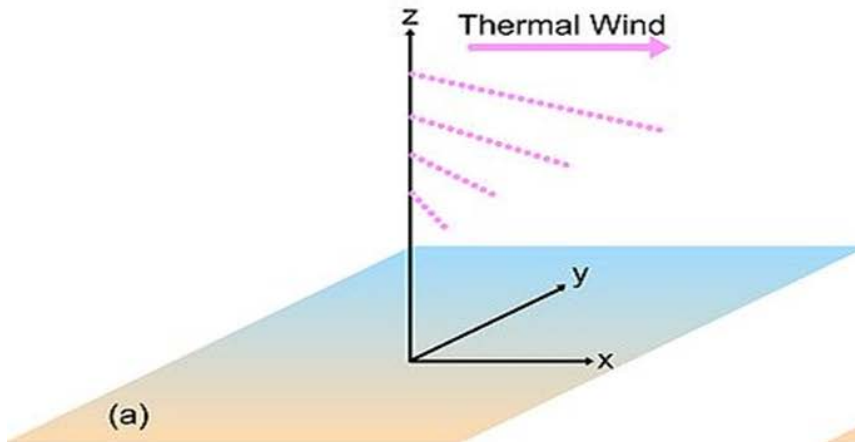
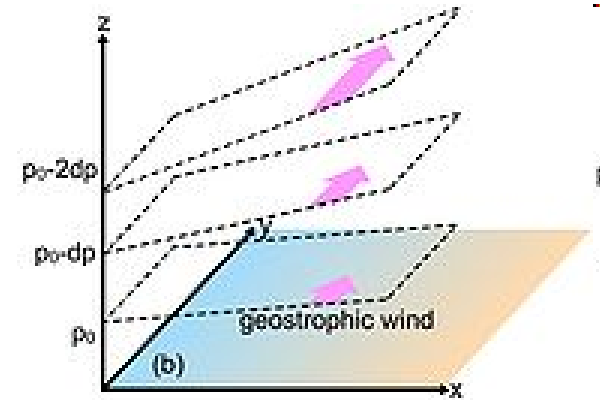
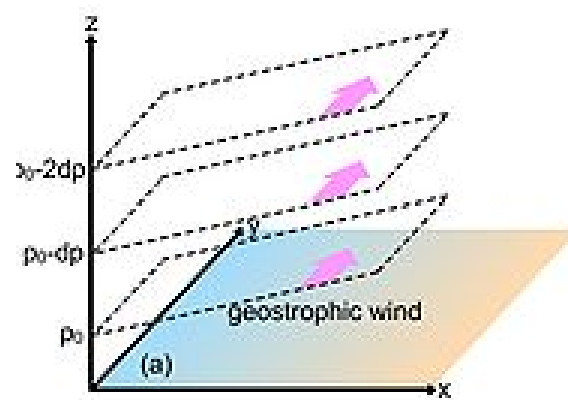
Thermal wind balance

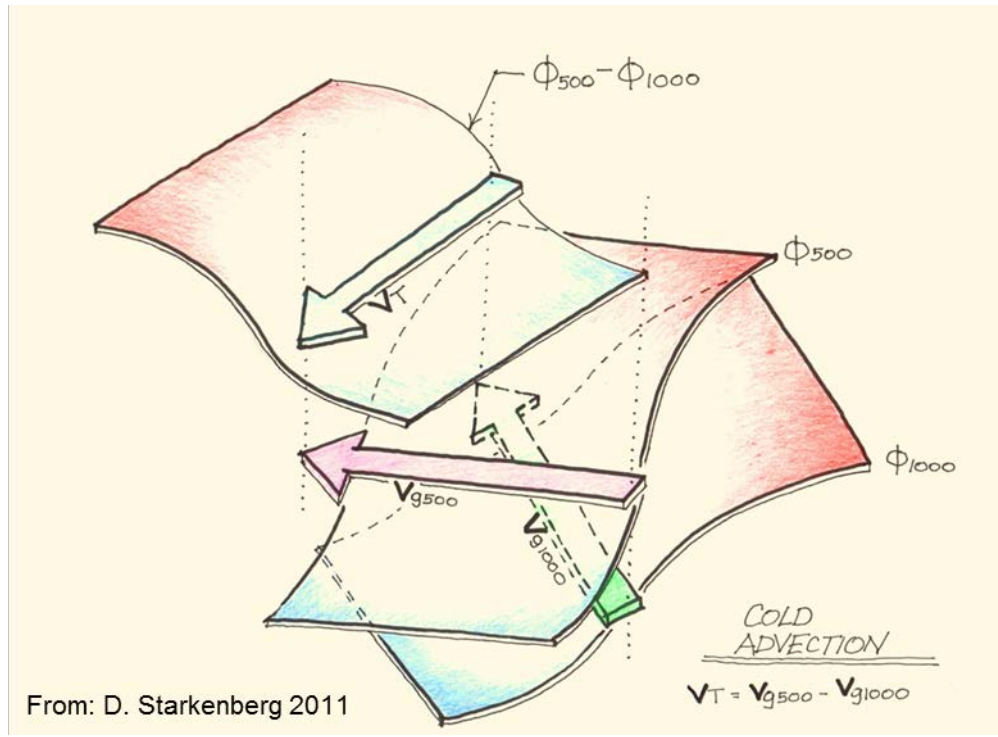
$$fv = \frac{R_d T}{p} \frac{\partial p}{\partial x} = R_d T \frac{\partial \ln p}{\partial x}$$

$$-\frac{g}{R_d T} = \frac{\partial \ln p}{\partial z}$$

$$f \frac{\partial u}{\partial z} \approx -\frac{g}{T} \frac{\partial T}{\partial y}$$

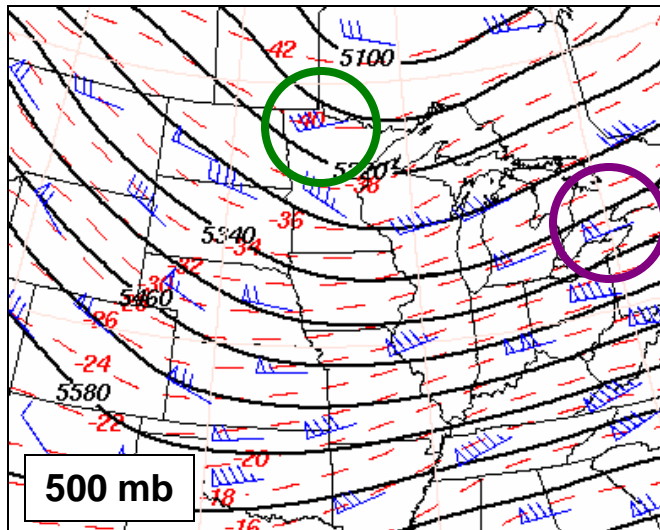
$$f \frac{\partial v}{\partial z} \approx \frac{g}{T} \frac{\partial T}{\partial x}$$



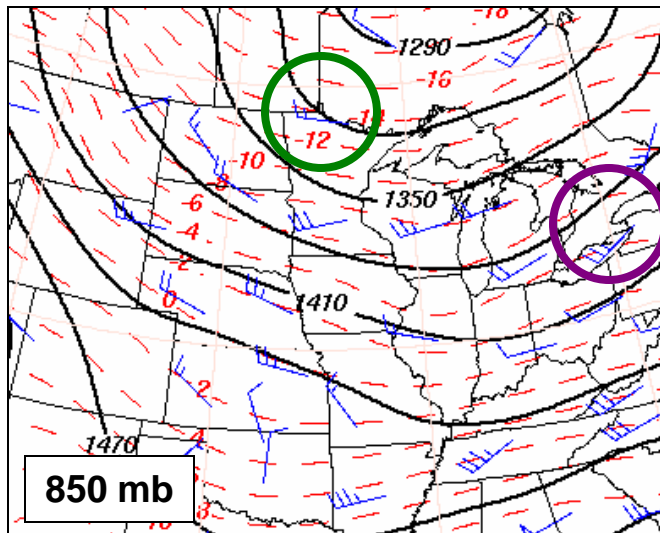


From: D. Starckenberg 2011

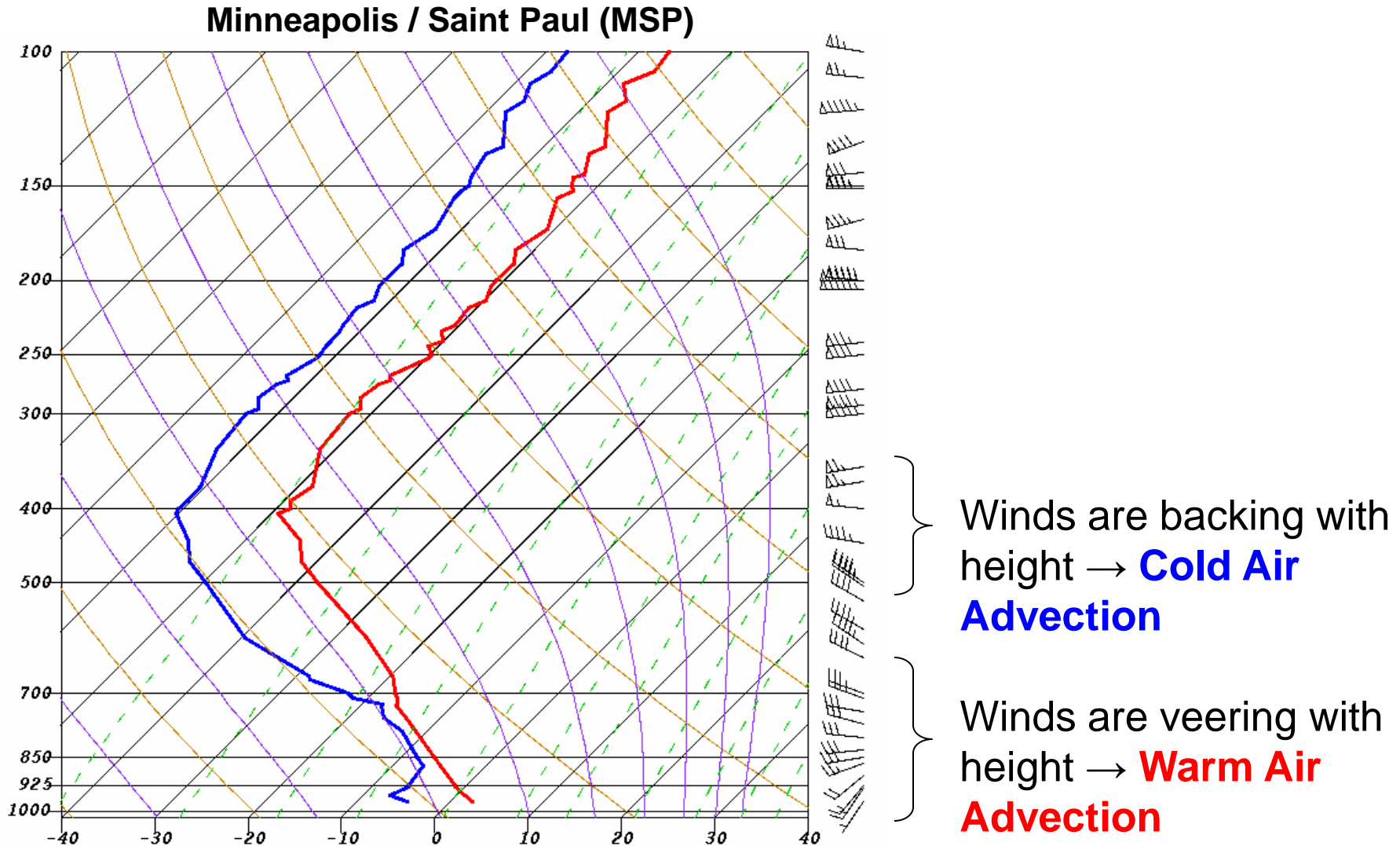
Thermal wind in synoptic application of pressure maps



International Falls, MN
Buffalo, NY



Thermal wind in synoptic application of soundings



Continuity equation

$$\frac{1}{\rho} \frac{\partial \rho}{\partial t} = \frac{\partial(\rho u)}{\partial x} + \frac{\partial(\rho v)}{\partial y} + \frac{\partial(\rho w)}{\partial z}$$

mass divergence form

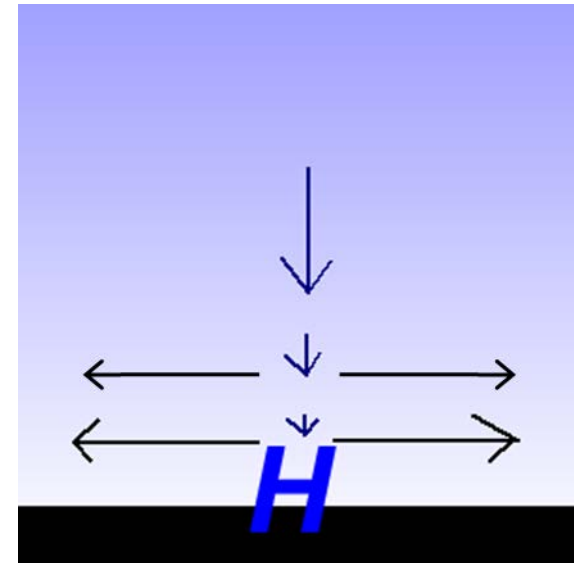
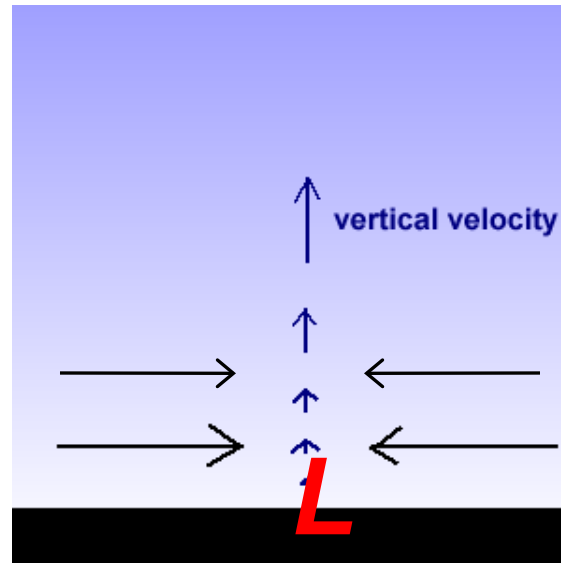
$$-\frac{1}{\rho} \frac{D\rho}{Dt} = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z}$$

velocity divergence form

$$0 = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z}$$

$$\frac{\partial w}{\partial z} = -\left[\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right]$$

$$\frac{\partial w}{\partial z} = -[\nabla \cdot \mathbf{V}_h]$$



Primitive equations

$$\frac{\partial(\rho \mathbf{v})}{\partial t} + \nabla \cdot (\rho \mathbf{v} \mathbf{v} + \mathbf{J}) = -\nabla p - \rho \nabla \Phi - 2\rho \boldsymbol{\Omega} \times \mathbf{v}$$

Momentum equation
(Newton's 2nd law)

local change

Coriolis force

Energy equation (1st law of thermodynamics;
conservation of energy)

$$\frac{\partial(\rho \Theta)}{\partial t} + \nabla \cdot (\rho \mathbf{v} \Theta + \mathbf{J}_{\Theta}) = S_{\Theta}$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

advection

Continuity equation
(conservation of mass)

$$\frac{\partial(\rho q_n)}{\partial t} + \nabla \cdot (\rho \mathbf{v} q_n + \mathbf{J}_{q_n}) = S_{q_n} \quad n = 0, \dots, 3$$

sources/sinks
(phase transitions)

Balance equations for dry air
and water substances

$$\frac{\partial(\rho \chi_m)}{\partial t} + \nabla \cdot (\rho \mathbf{v} \chi_m + \mathbf{J}_{\chi_m}) = S_{\chi_m} \quad m = 1, \dots, M$$

Balance equations for M
species and particles

sources/sinks
(reactions, photolysis)

Equation of state

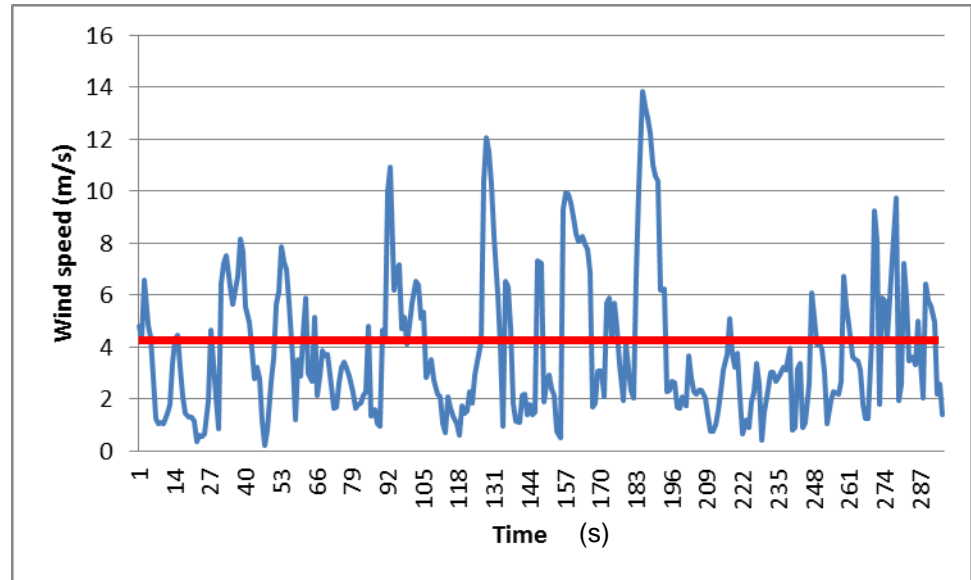
$$p = \rho R_0 T_v$$

Mean and perturbation plus Reynolds-averaging

$$p = \bar{p} + p'$$

$$u = \bar{u} + u'$$

$$\rho = \bar{\rho} + \rho'$$



Mean and perturbation plus Reynolds-averaging

$$\overline{\varphi + \gamma} = \overline{\varphi} + \overline{\gamma}$$

$$\overline{\varphi\gamma} = \overline{\varphi\gamma} + \overline{\varphi'\gamma'}$$

$$\overline{\alpha\varphi} = \alpha\overline{\varphi} \text{ with } \alpha = \text{const.}$$

$$\overline{\overline{\varphi}} = \overline{\varphi}$$

$$\overline{\frac{\partial\varphi}{\partial t}} = \frac{\partial\overline{\varphi}}{\partial t}$$

$$\overline{\nabla \cdot \varphi} = \nabla \cdot \overline{\varphi}$$

$$\overline{\nabla\varphi} = \nabla\overline{\varphi}$$

$$\overline{\varphi'} = 0$$

Example: Boussinesq approximation

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

$$\frac{\partial(\rho \mathbf{v})}{\partial t} + \nabla \cdot (\rho \mathbf{v} \mathbf{v} + \mathbf{J}) = -\nabla p - \rho \nabla \Phi - 2\rho \boldsymbol{\Omega} \times \mathbf{v}$$

$$\frac{\partial \bar{\rho}}{\partial t} + \nabla \cdot (\bar{\rho} \bar{\mathbf{v}} + \underbrace{\bar{\rho}' \bar{\mathbf{v}}'}_{\text{neglected}}) = \frac{\partial \bar{\rho}}{\partial t} + \nabla \cdot (\bar{\rho} \bar{\mathbf{v}}) = 0$$

$$\begin{aligned} \frac{\partial(\bar{\rho} \bar{\mathbf{v}} + \underbrace{\bar{\rho}' \bar{\mathbf{v}}'}_{\text{neglected}})}{\partial t} + \nabla \cdot \left(\bar{\rho} \bar{\mathbf{v}} \bar{\mathbf{v}} + \bar{\rho} \bar{\mathbf{v}}' \bar{\mathbf{v}}' + \underbrace{\bar{\rho}' \bar{\mathbf{v}}' \bar{\mathbf{v}} + \bar{\mathbf{v}} \bar{\rho}' \bar{\mathbf{v}}' + \bar{\rho}' \bar{\mathbf{v}}' \bar{\mathbf{v}}'}_{\text{neglected}} + \bar{\mathbf{J}} \right) \\ = -\nabla \bar{p} - \bar{\rho} \nabla \phi - 2\bar{\rho} (\boldsymbol{\Omega} \times \bar{\mathbf{v}}) - 2\boldsymbol{\Omega} \times \underbrace{(\bar{\rho}' \bar{\mathbf{v}}')}_{\text{neglected}} \end{aligned}$$

identity tensor

Stokes friction tensor

Reynolds stress tensor

$$\frac{d\bar{\rho}}{dt} = -\bar{\rho} \nabla \cdot \bar{\mathbf{v}}$$

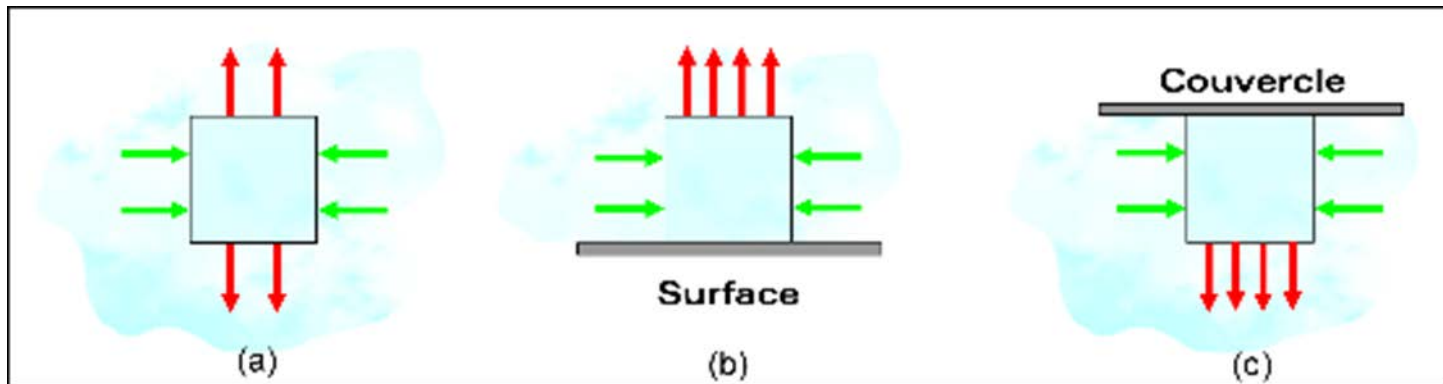
$$\bar{\rho} \frac{d\bar{\mathbf{v}}}{dt} + \nabla \cdot (\bar{\rho} \mathbf{E} + \bar{\mathbf{J}} + \mathbf{F}) = -\bar{\rho} \nabla \phi - 2\bar{\rho} \boldsymbol{\Omega} \times \bar{\mathbf{v}}$$

Anelastic approximation

$$\frac{d\rho}{dt} = -\rho \nabla \mathbf{u} \quad \Rightarrow \quad \nabla(\rho_{ref} \mathbf{u}) = 0$$

ρ fully prognostic

→ ρ diagnostic (ideal gas law)



Closure problem – 1st order closure

$$\mathbf{F} = \bar{\rho} \overline{\mathbf{v}' \mathbf{v}'} = -\bar{\rho} \mathbf{K}_m : \left(\nabla \bar{\mathbf{v}} + (\nabla \bar{\mathbf{v}})^T \right)$$

$$\mathbf{H} = c_{p,0} \bar{\rho} \overline{\mathbf{v}' \Theta'} = -c_{p,0} \bar{\rho} \mathbf{K}_h \cdot \nabla \bar{\Theta}$$

$$\mathbf{F}_{q_n} = \bar{\rho} \overline{\mathbf{v}' q_n'} = -\bar{\rho} \mathbf{K}_{q_n} \cdot \nabla \bar{q}_n \quad \text{for } n = 1, 2, 3$$

$$\mathbf{F}_{\chi_m} = \bar{\rho} \overline{\mathbf{v}' \chi_m'} = -\bar{\rho} \mathbf{K}_{\chi_m} \cdot \nabla \bar{\chi}_m \quad \text{for } m = 1, \dots, M.$$

References

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